

Spin dynamics in Microcavities



M. Vladimirova
S. Cronenberger
D. Scalbert
PhD students:

A. Brunetti (2004-2007)
R. Giri (2010□)

**Laboratoire Charles Coulomb, CNRS, Université
Montpellier 2, France**

A. Miard
A. Lemaître
J. Bloch

**Laboratoire de Photonique et de Nanostructures,
CNRS, Marcoussis, France**

R. André

**Institut Neel, CNRS, Université Joseph Fourier
Grenoble, France**

A. V. Kavokin

**Physics and Astronomy School, University of
Southampton, UK**

K. V. Kavokin

A. F. Ioffe Institute, St-Petersburg, Russia

G. Malpuech
D. Solnyshkov

LASMEA, Clermont-Ferrand, France

M. Nawrocki

Warsaw University, Poland

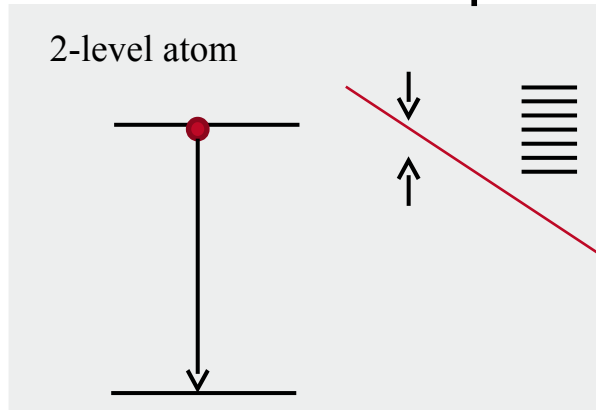


Outline

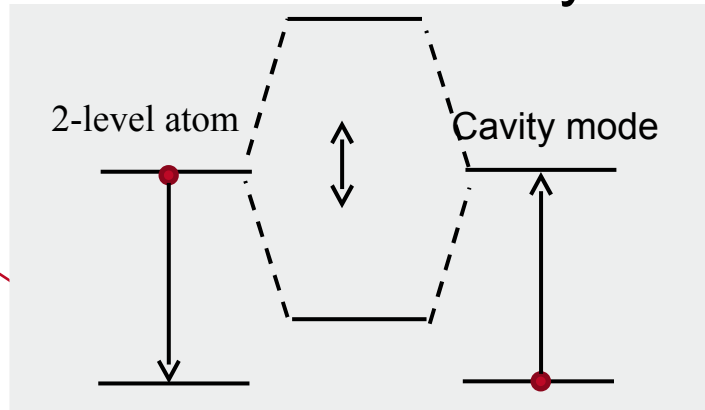


Vacuum field Rabi oscillation

Atom in free space

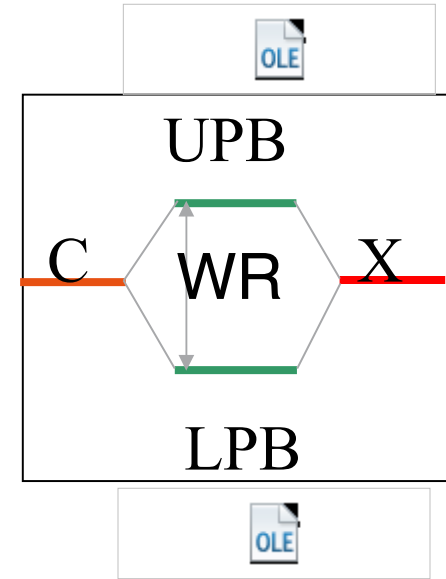
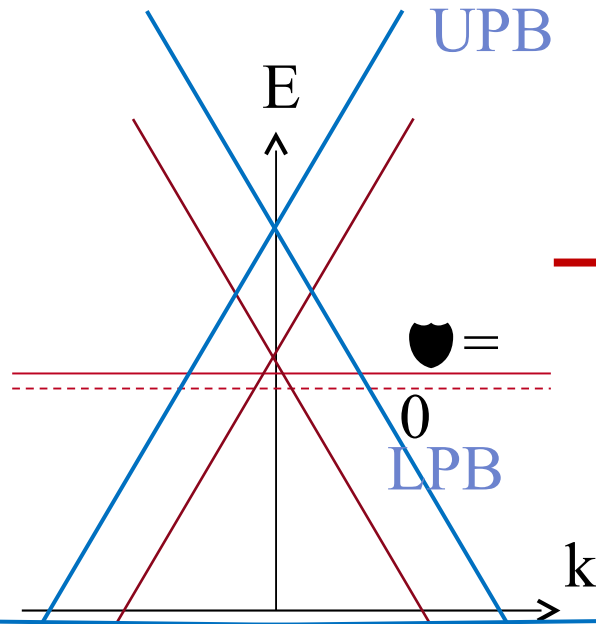
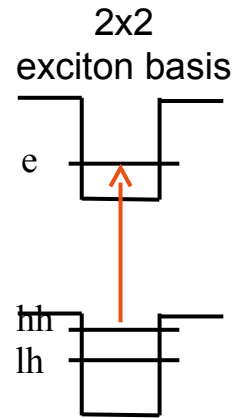
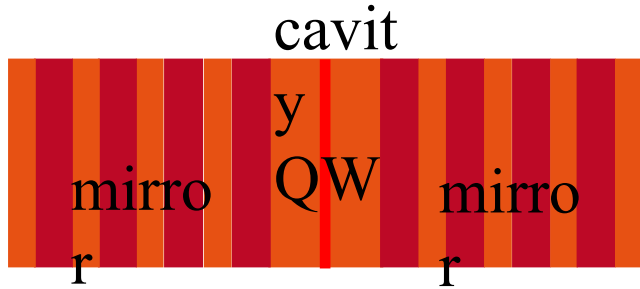


Atom in cavity

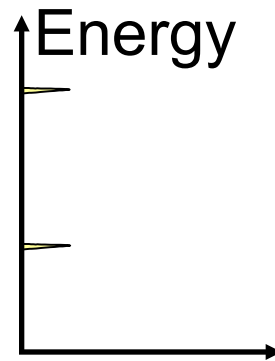
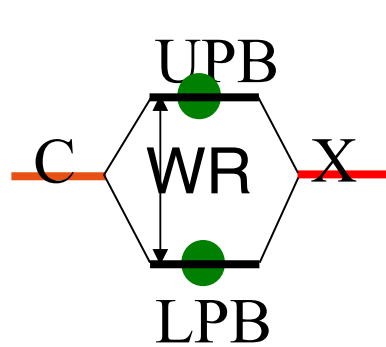
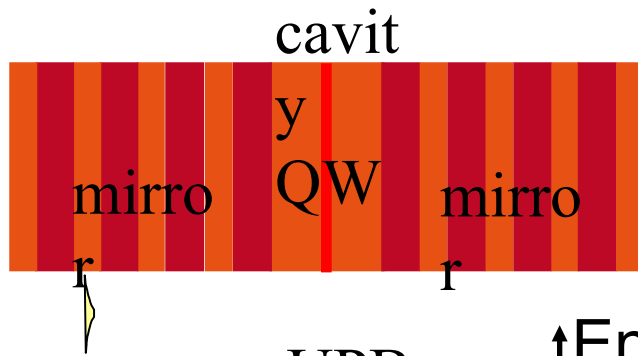


Spontaneous emission Vacuum field Rabi oscillation

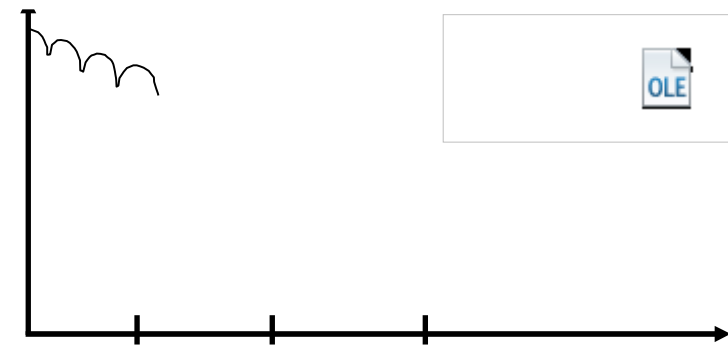
Microcavity in the strong coupling regime



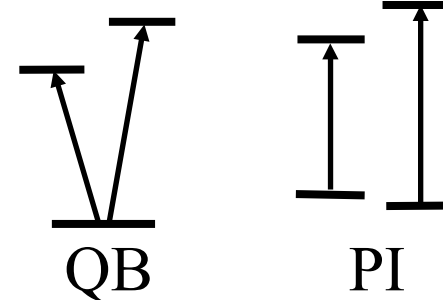
QW microcavity : time- resolved transmission of the light



Linear transmission



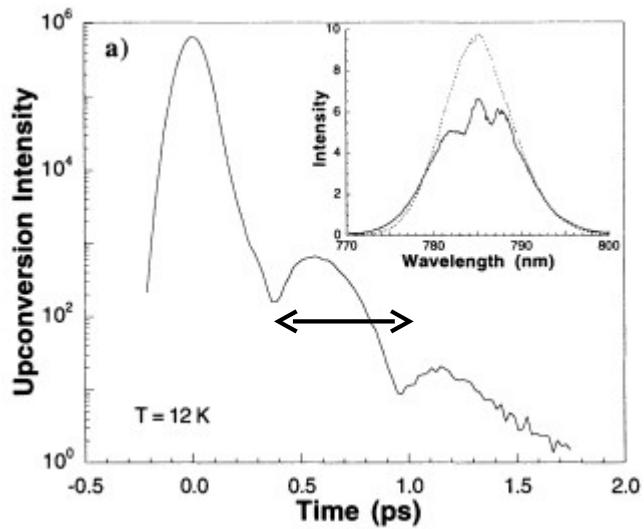
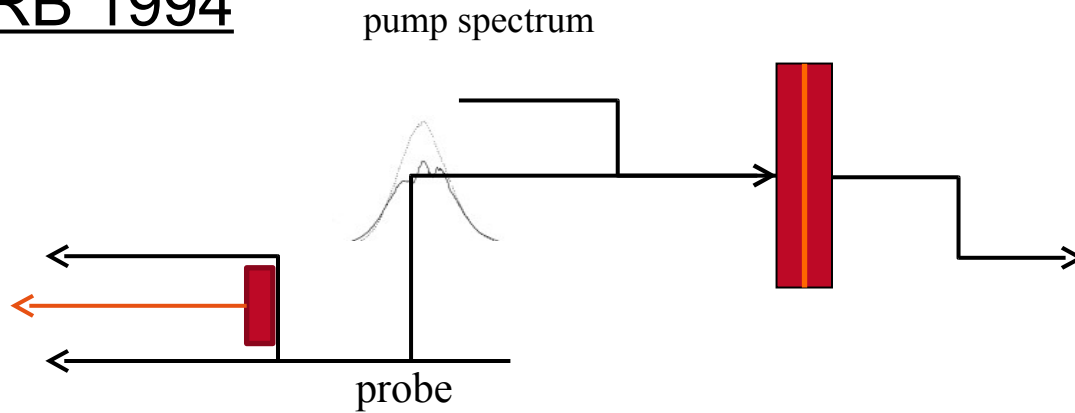
Transmission 0 $\frac{p \square \square}{WR}$ Time



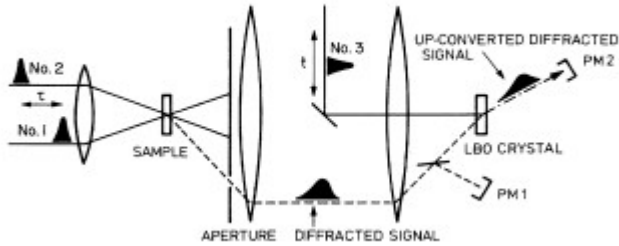
These oscillations can result either from polarization interference (PI) in the detector, or from quantum beats (QB) but **cannot be distinguished in linear optics**

QW microcavity : time- resolved emission

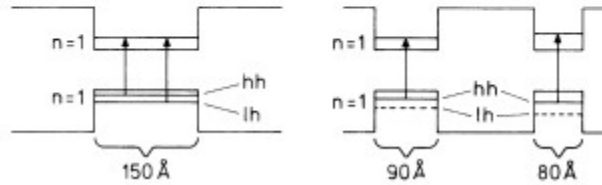
Norris et al, PRB 1994



Distinction between QB and PI by time-resolved FWM

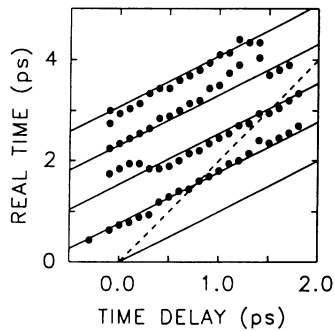


M. Koch et al, PRB (1993)
Two-pulse photon echo
experiment

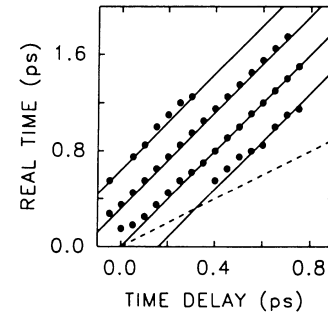


QB

PI



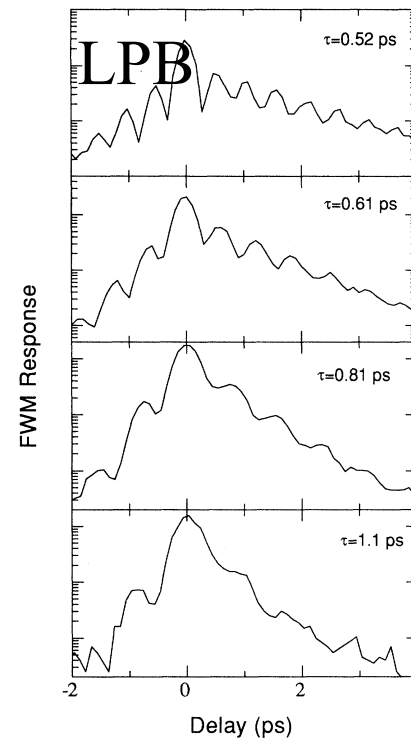
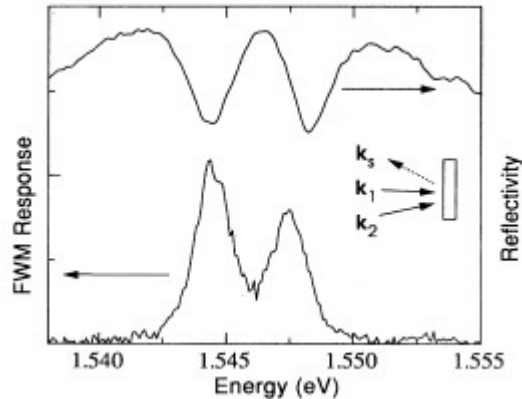
maxima at
 $t=t+nTB$



maxima
at
 $t=2t+nT$
B

M. Koch et al, PRL **69**, 3631 (1992)

Spectrally resolved FWM: GaAs-based microcavity

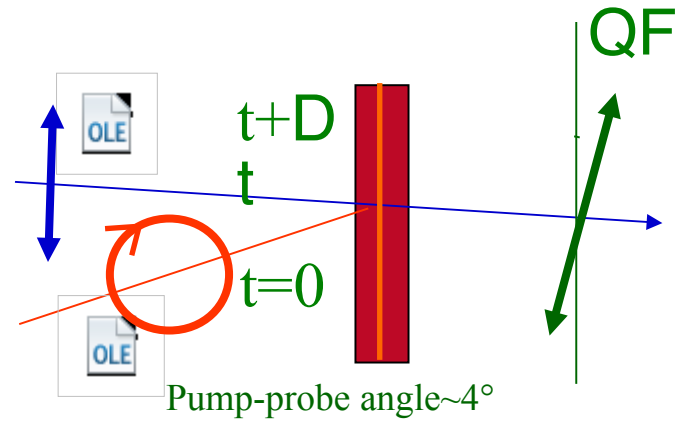
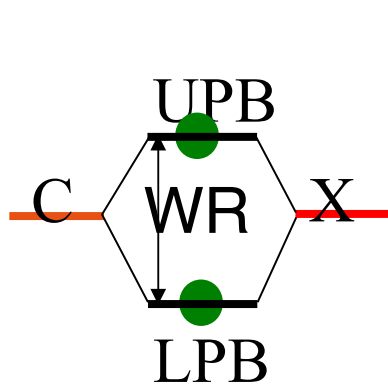


Spectral filtering of the signal prevents PI effects ☾ QBs are observed

QBs are observed even at large detunings: nonlinear response is due to excitonic part of the wavefunction

Wang et al, PR B **51**, 14713 (1994)

Time-resolved Faraday rotation



Samples

1/2 cavity



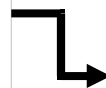
GaAs/GaAlAs

CdMnTe/CdMgTe



GaInAs QW

CdMnTe QW



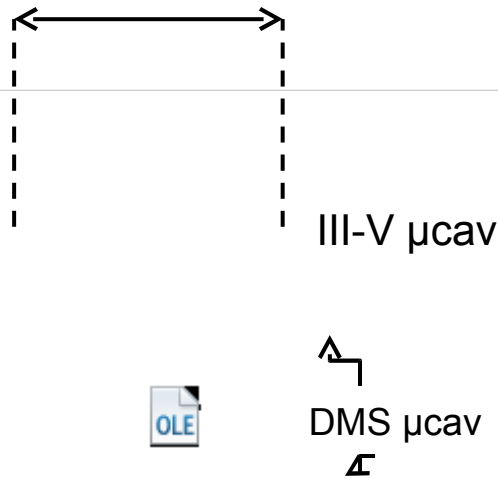
sample	QW	Lw (nm)	WR (meV)	Q	tc (ps)	TR (ps)
M1025	CdMn5%Te	7	6	400	0.25	0.6
M992	CdMn0.7%Te	6	10	400	0.25	0.4
11G20	GaIn5%As	8	3.5	2500	1.8	1.1

GaInAs MC : J. Bloch, LPN (Paris)

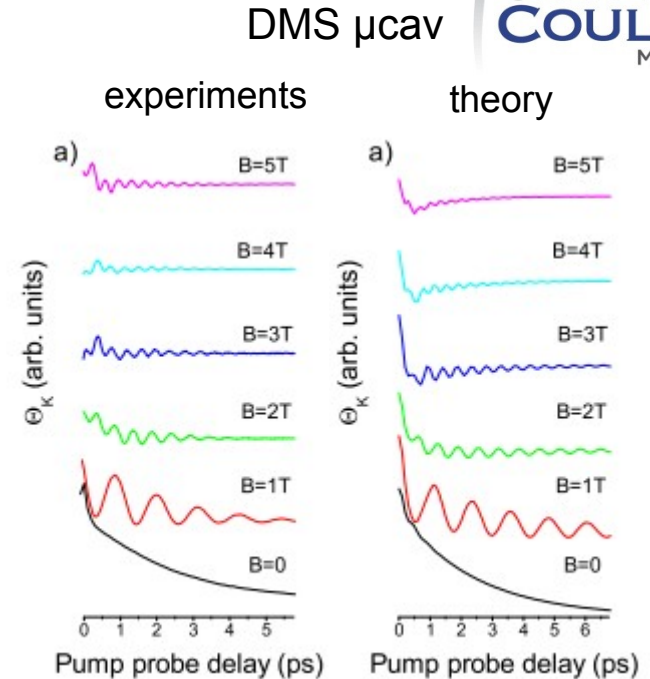
CdMnTe MC : R. André, LSP (Grenoble)



Spectrally integrated TRFR



Scalbert et al Proc. PLMCN7 2008

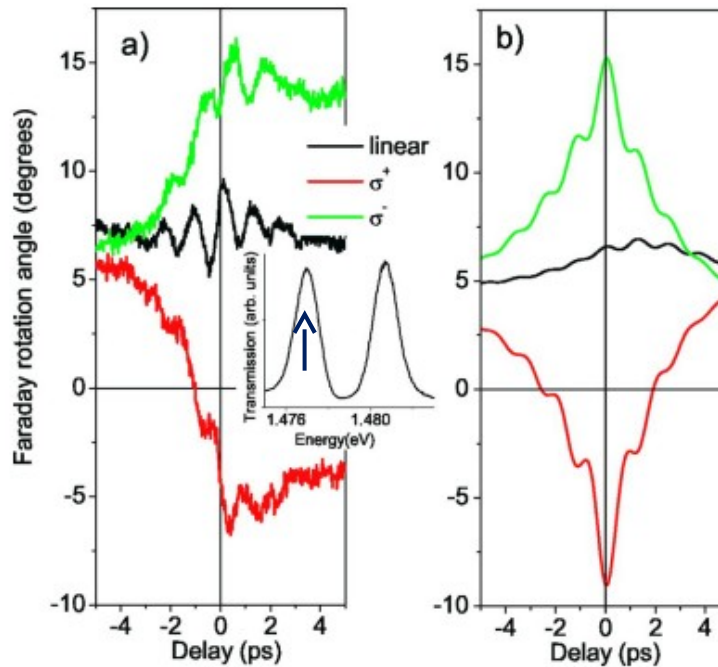


Brunetti et al PRB 2006

- Rabi oscillations seen only in the μ cavity with long enough cavity lifetime
- Oscillations better seen at negative delays
- Long living non oscillating decay probably due to spin polarized excitons from the reservoir

• PI effects not seen in Kerr rotation: but do we really see Rabi oscillations?

Spectrally-resolved photo-induced Faraday Rotation

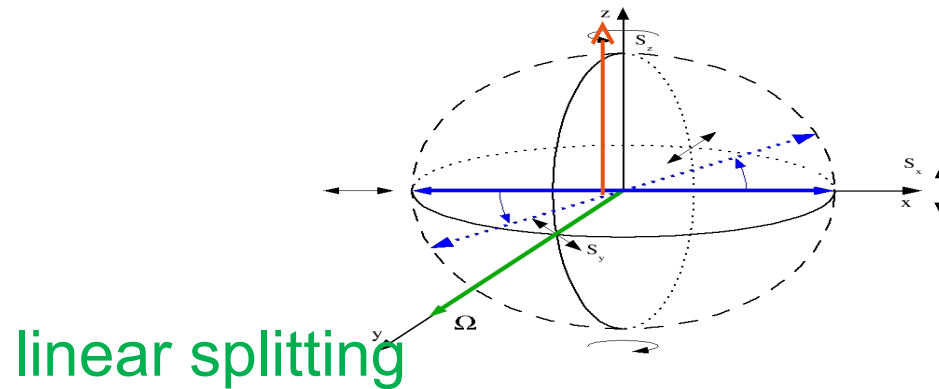


- Faraday rotation exhibits beatings with Rabi period 1.25 ps
- Linear birefringence induces
 - rotation of probe polarization at negative delays
 - Existence of beatings for linear polarization of the pump

Brunetti et al, PRB 2006

Conversion from linear to circular polariton state due to linear splittings of polariton branches: Poincaré sphere representation

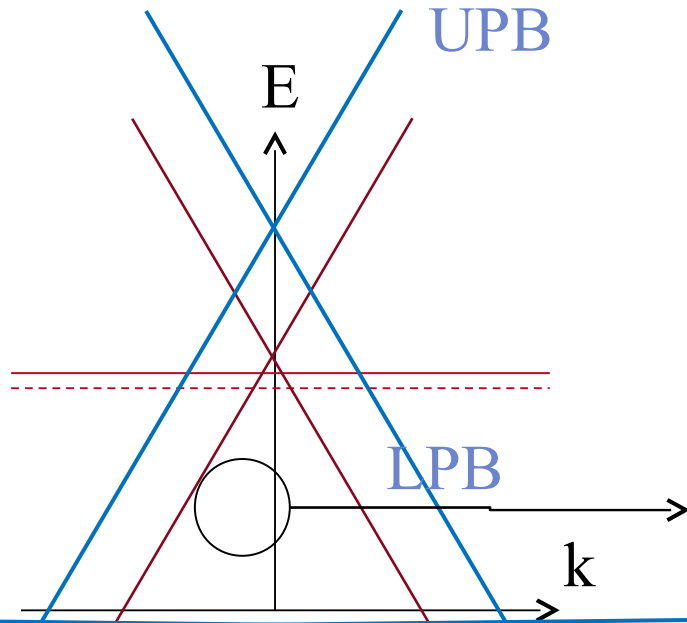
left-right splitting



See review by I. Shelykh, A. Kavokin, G. Malpuech PSS(b) 2005

Spin state of the condensate

1) Ideal system (no splitting)



TE-TM
splitting

$k=0$

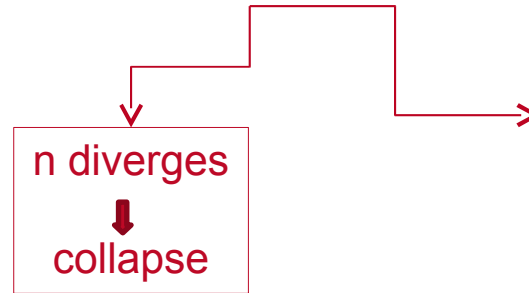
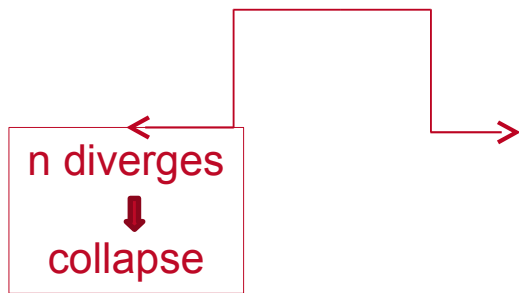
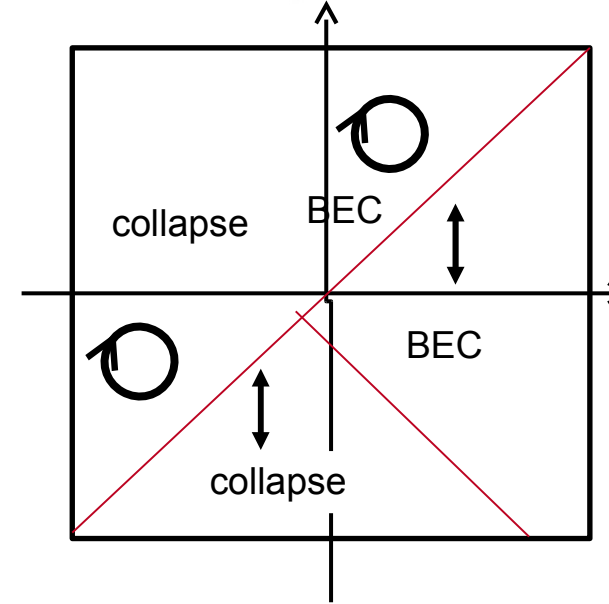
- Fundamental state is degenerate \hookrightarrow condensate must select a specific state among an infinity of states
- Spontaneous symmetry breaking with build-up of stochastic linear polarization is expected (Shelykh et al

Semicond. Sci. Technol. 2010)

Free energy of the condensate



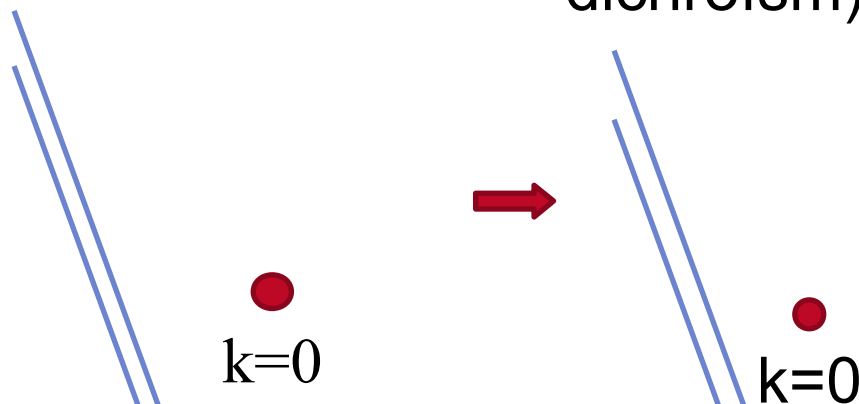
Minimize free energy



Spin state of the condensate

2) Real system

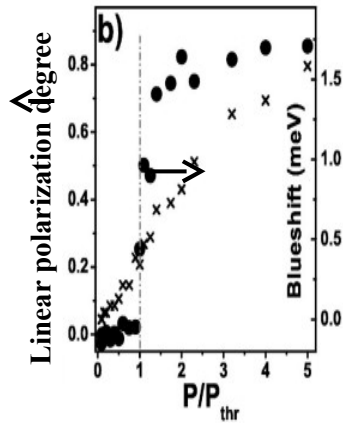
- a) asymmetric QW }
b) strain in the mirrors } \hookrightarrow lower symmetry (linear dichroism)



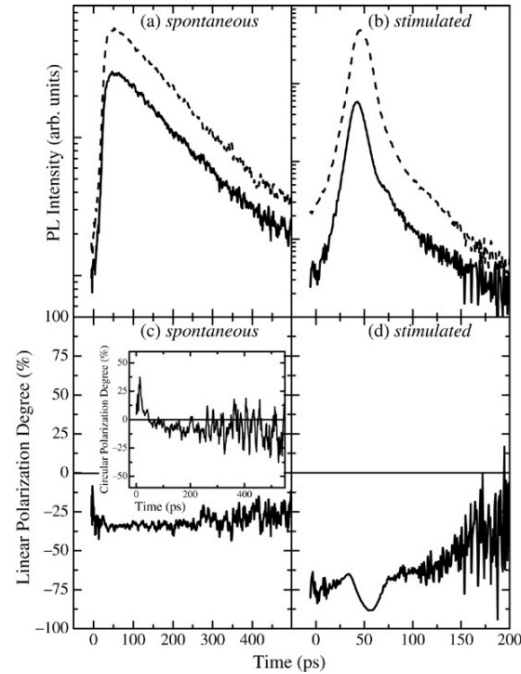
Fundamental state: degeneracy is removed \hookrightarrow
polarization of the condensate is pinned to a fixed
axis

see Shelykh et al Superlattices and Microstructures (2007)

Pinning of the polarization of light emitted by a microcavity



Kasprzak et al PRB 2007



Klopotowski *et al.*, SSC 2006

CdTe/CdMgTe μ cav

Non resonant excitation

Polarization fixed with respect to cristal axis

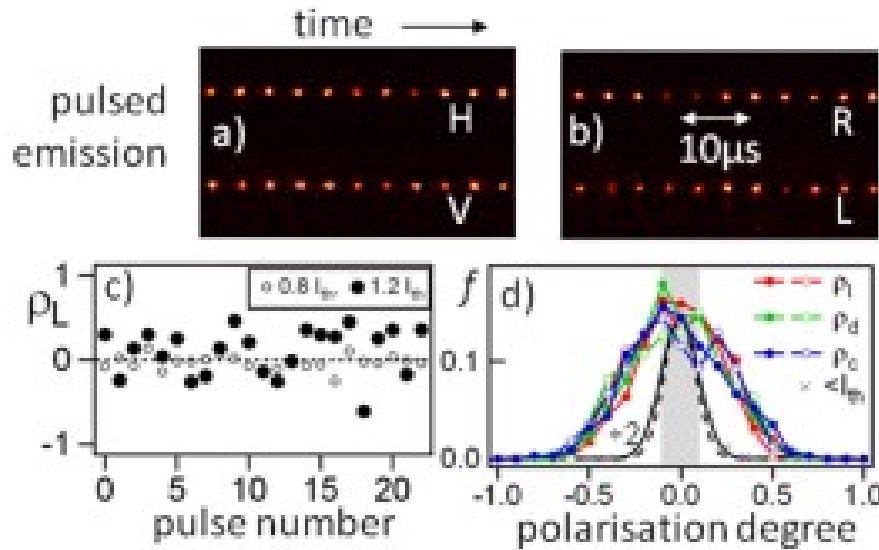
increases above threshold for stimulated emission

See also:

Kasprzak et al Nature 2006

Balili et al Science 2007

Spontaneous symmetry breaking in a BEC



Polarization resolved
emission above threshold

Polarization histogram

Bulk GaN μcav :

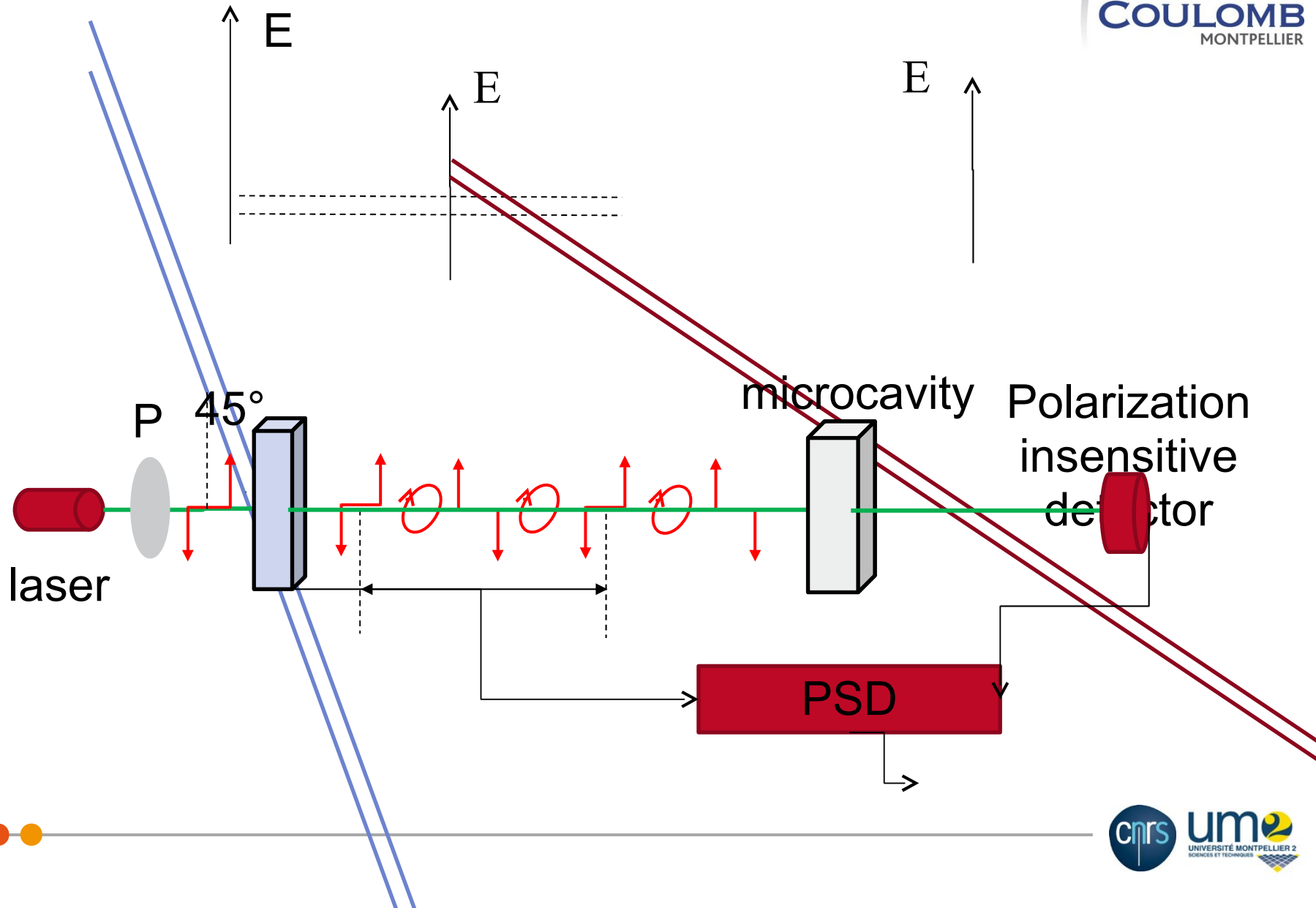
No strain-induced birefringence

spin-isotropic polariton-polariton interactions

Room-temperature

Baumberg et al, PRL 2008

Polarization resolved transmission



Linear dichroism: results

LPB

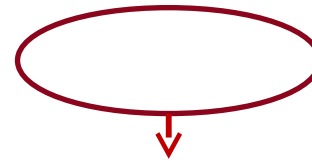
UPB



OLE

(11G20)

Origin of the anomalous signal?



Non-linear optical effect

3rd order nonlinear polarization

Boyd, Nonlinear optics



Power dependence of signal



Conclusion

UPB : linear optical effect / linear dichroïsm dominates

LPB : nonlinear optical effect / mixed dichroïsm dominates

How to separate the 2 contributions?

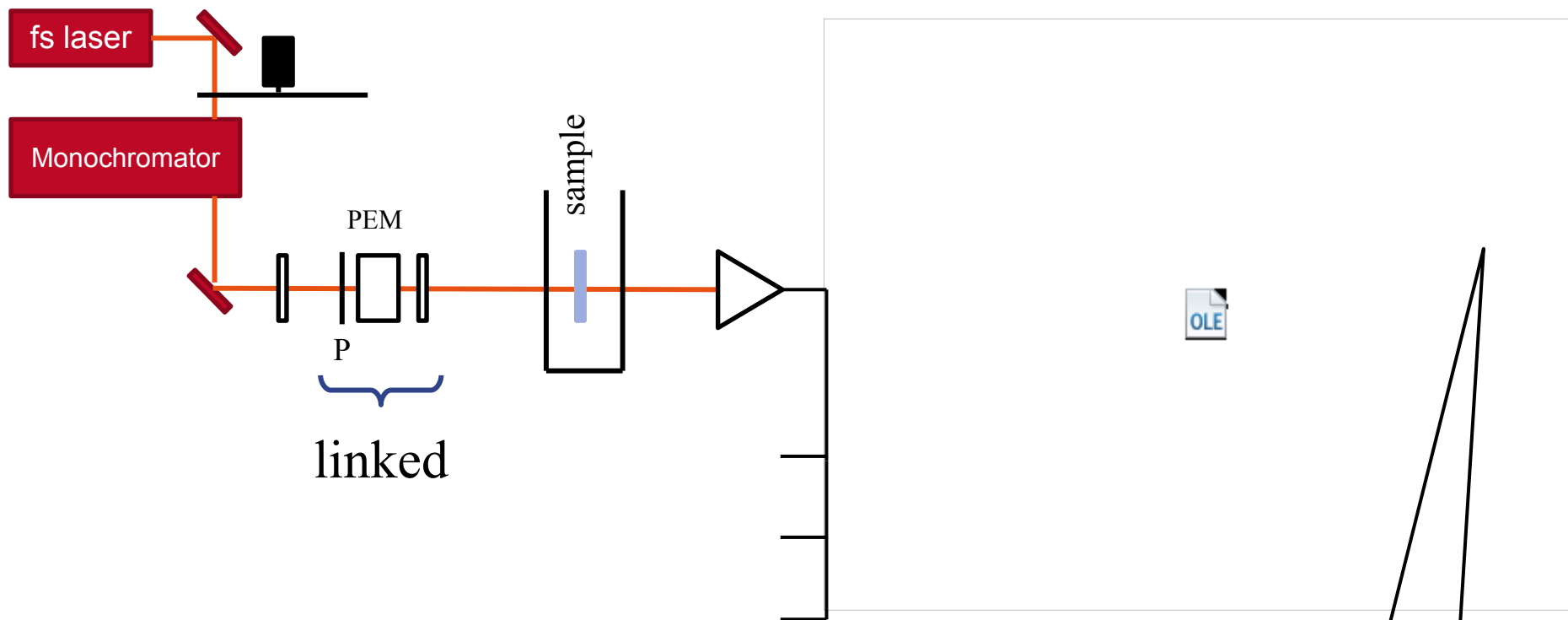
New modulation scheme



linear dichroïsm

mixed dichroïsm

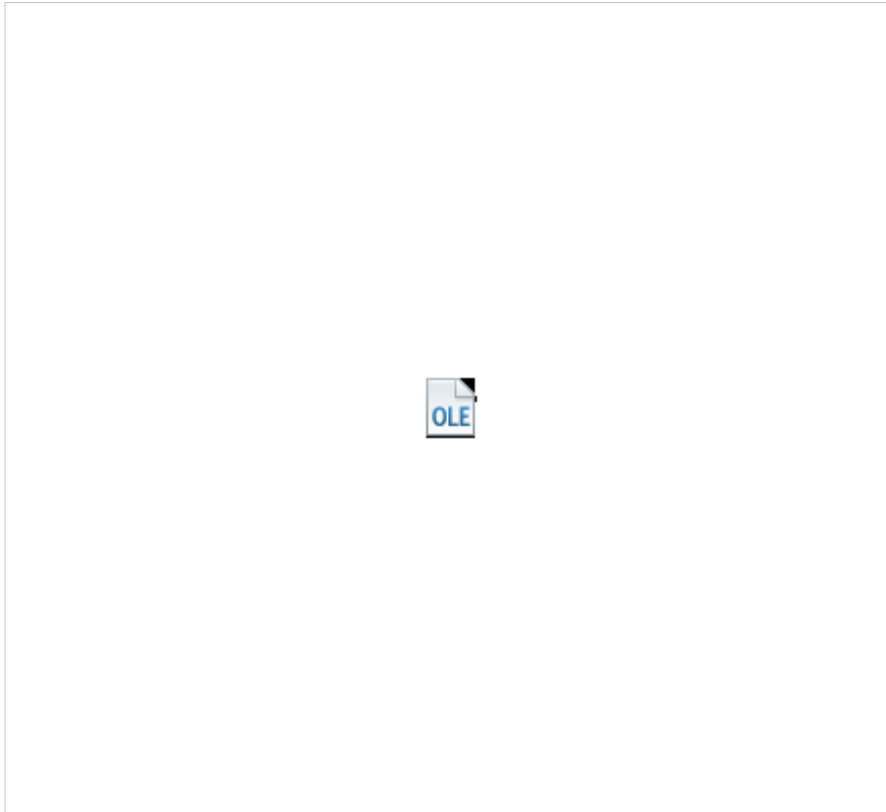
Simultaneous detection of linear and mixed dichroism



Linear and mixed dichroism do not spoil each other

Mixed dichroism also appears on UPB but with opposite sign

Lineshape of linear dichroism is different on LPB and UPB



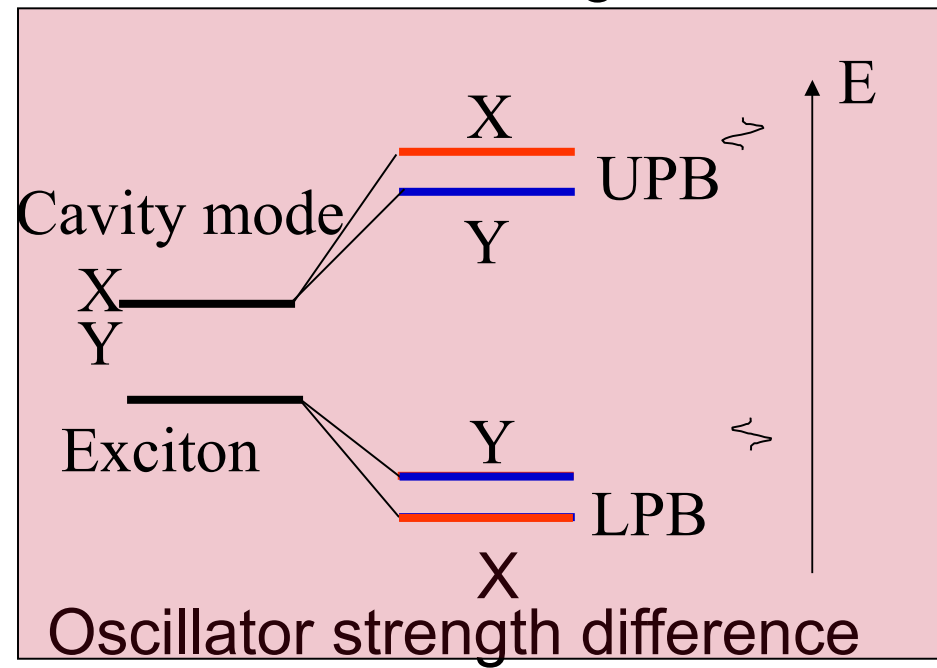
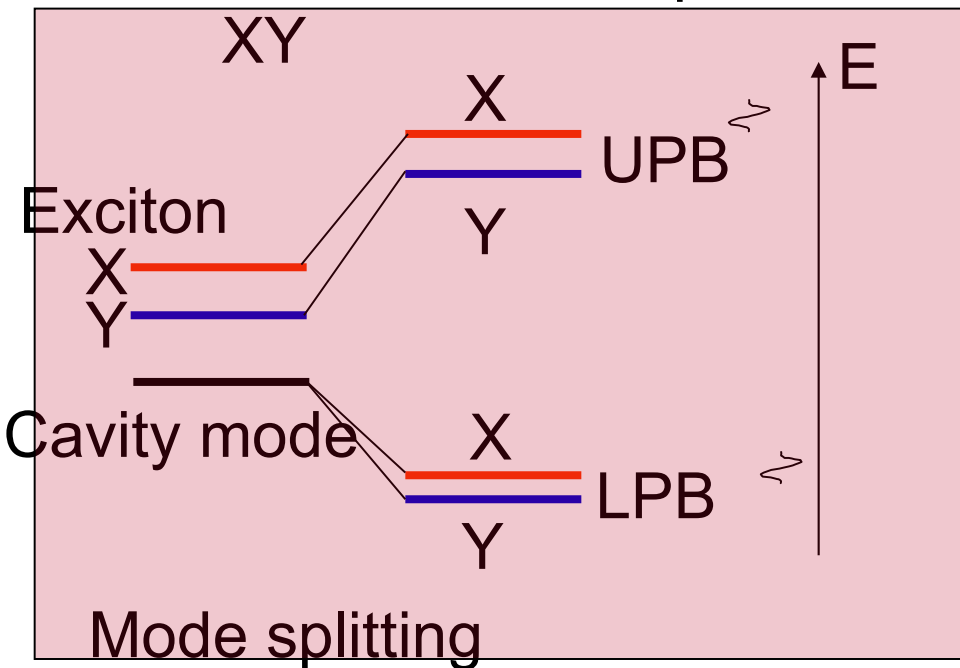
LD of UPB is stronger than LD of LPB

principal axis of LD are different for UPB and LPB

Linear dichroism : discussion

3 contributions to polariton linear splitting :

- Exciton splitting XY
- cavity mode splitting X'Y'
- Polarization dependent exciton oscillator strength



Linear dichroism : results and discussion

C

X

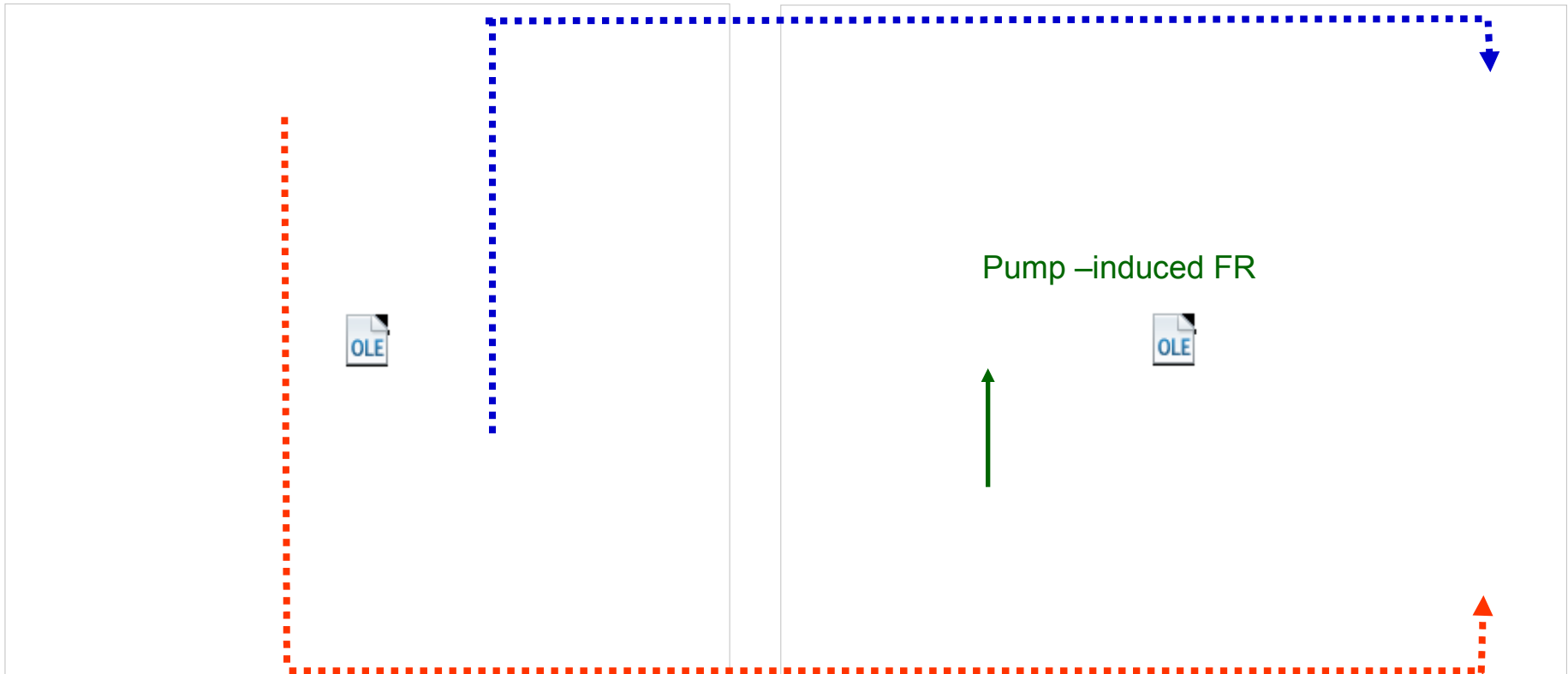


C

X

- Exciton splitting ~ 15 meV
- Cavity mode splitting ~ 20 - 30 meV
- Polarization dependent Rabi splitting ~ 20 meV
- $\sim 30^\circ$ between dichroism axes for exciton and cavity mode

• splittings are smaller than those observed in PL experiments
Brunetti et al. PLMCN6 (2006)
Scalbert et al. PLMCN7 (2007)

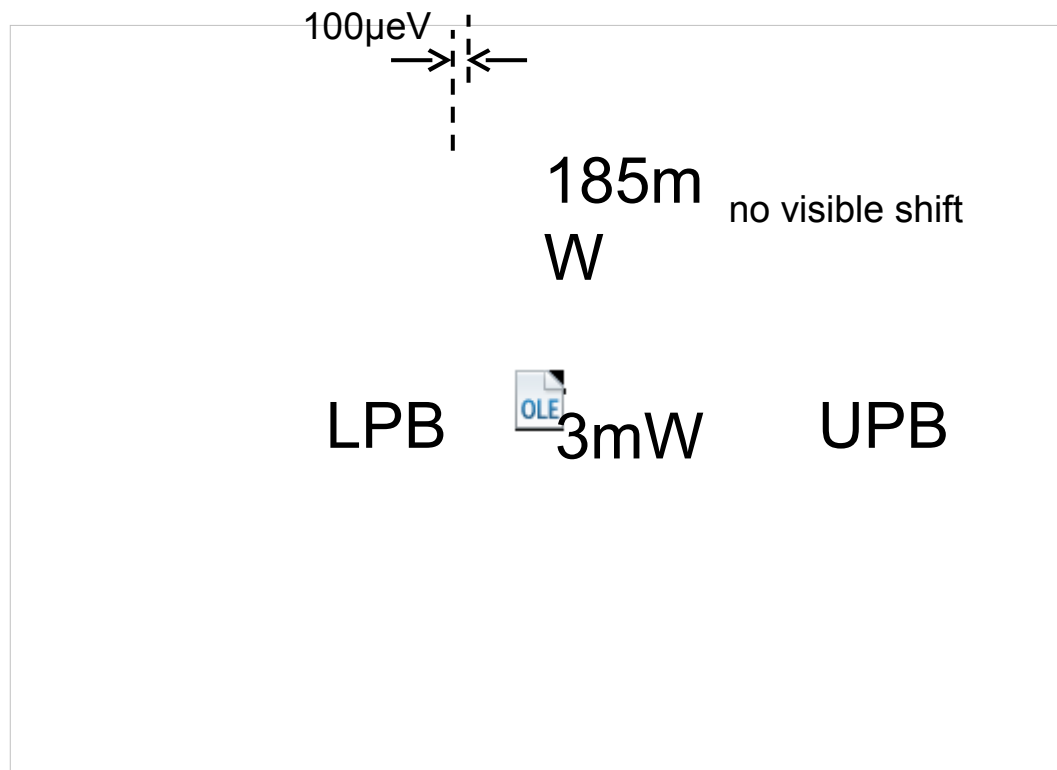


Transmission difference up to 40% $1 \text{ mW} \leftrightarrow 9 \cdot 10^8$

Much stronger at LPB than at UPB photons/cm²

: $T_c > T_l$

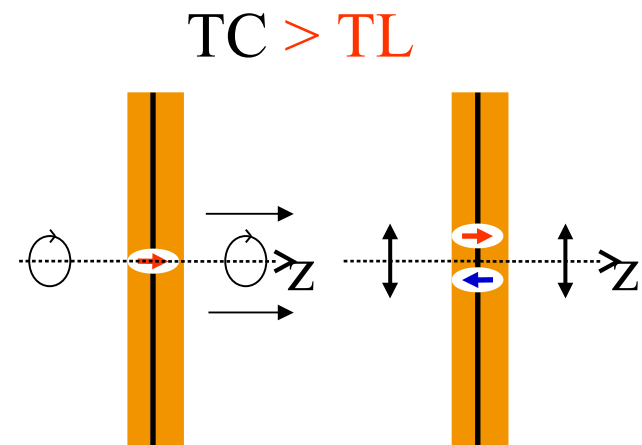
Polarization resolved transmission



Line broadening :6%

Splitting : 100 meV

Main effect : transmission
difference up to 40%



It may result from spin
dependent polariton-polariton
interactions

Microscopic origin of the optical nonlinearity?

dominant
↙ contribution

Phase-space filling

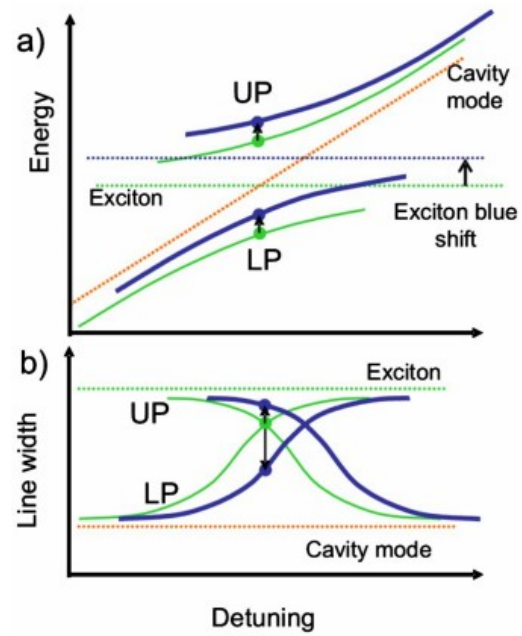
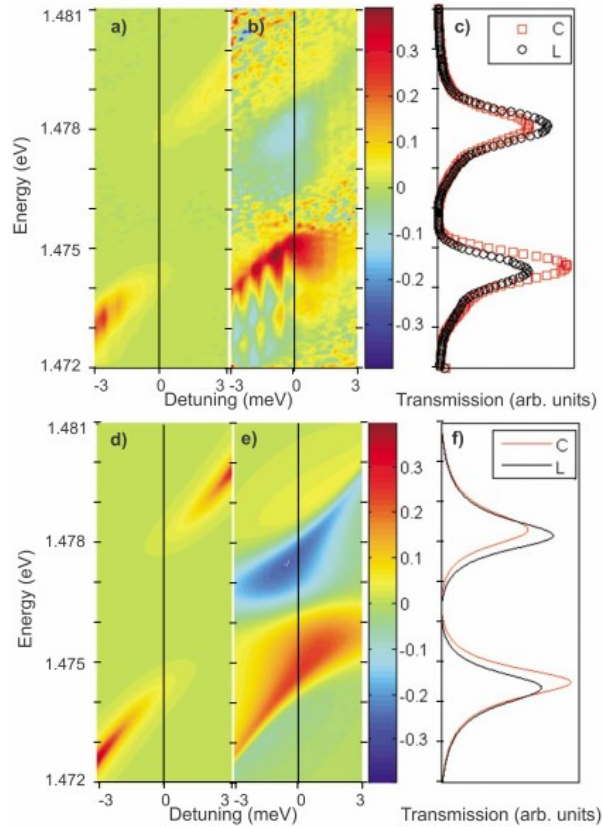
Exciton-exciton
interactions

In agreement with observed linear shifts

LPB

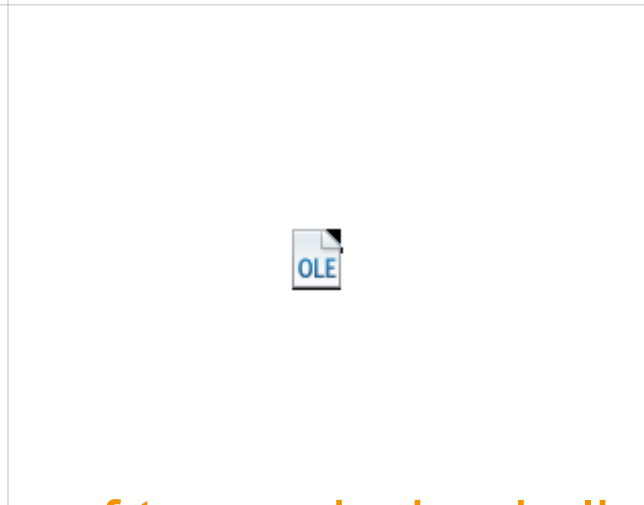
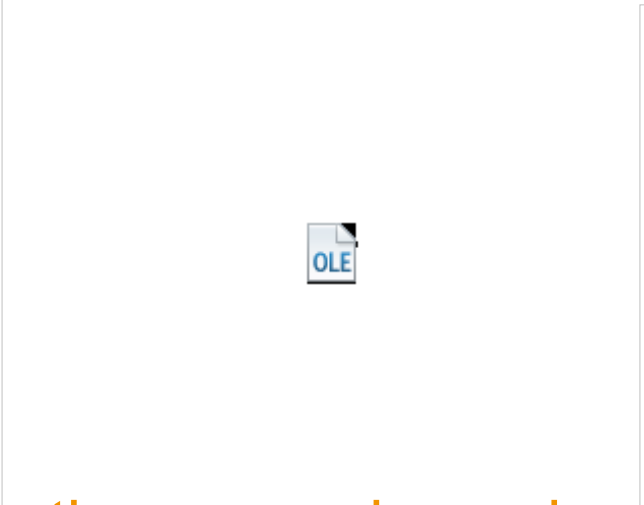
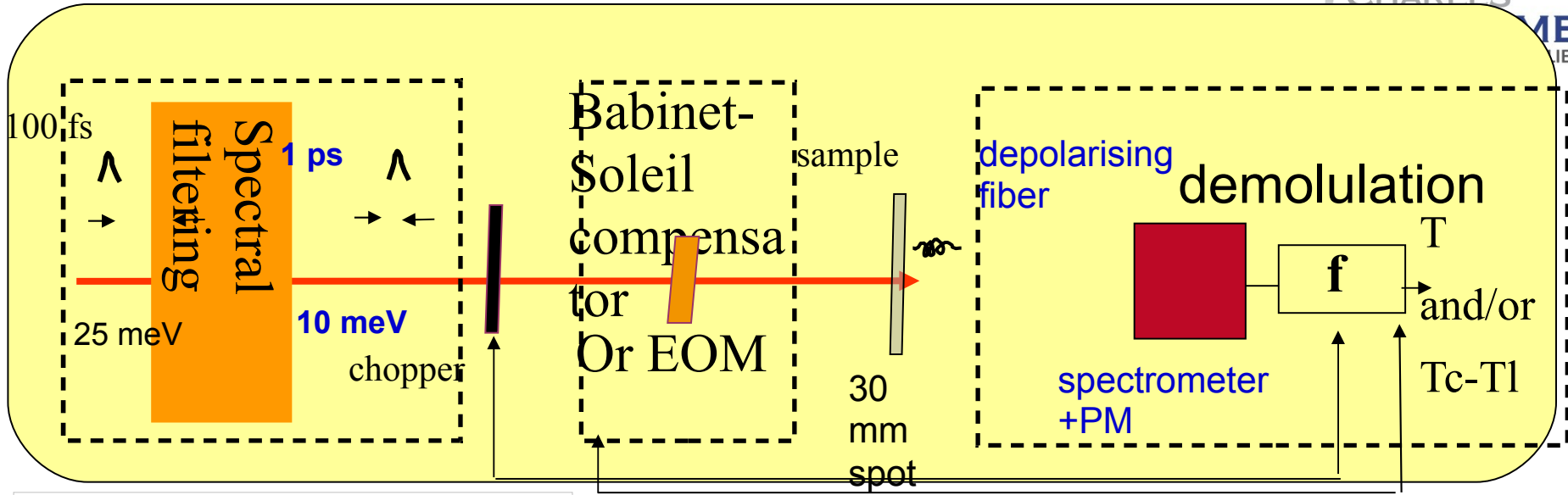


Interpretation of mixed dichroism: spin-dependent blue shift



Vladimirova et al, PRB 2009

Polariton energy shift from transmission experiments



GaAs 1/2 cavity,
In_{0.5}Ga_{0.95}As QW,

GaAs/Ga_{0.9}Al_{0.1}As
Bragg mirrors

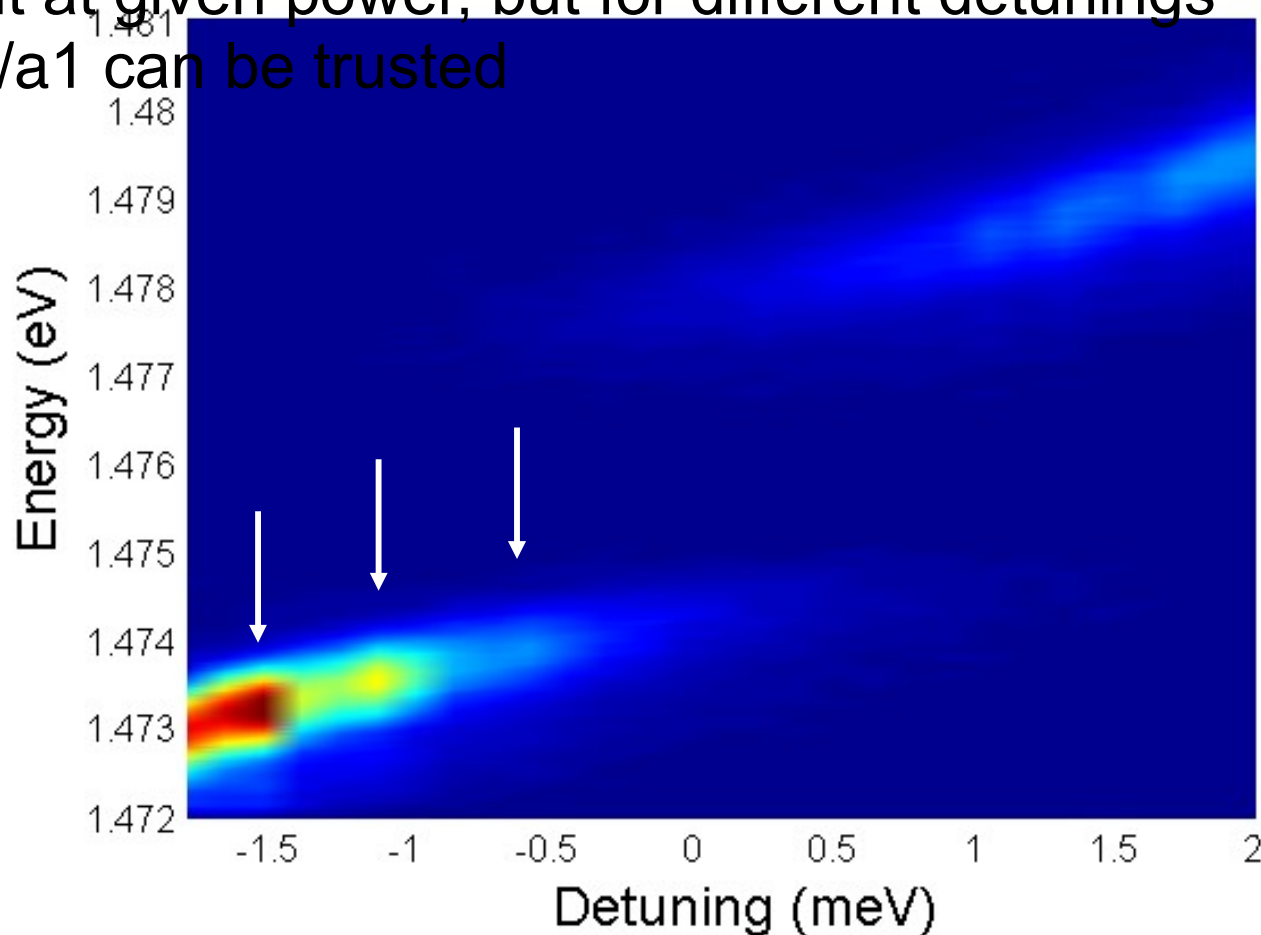
23 pairs/29pairs
WR=3.5 meV

for the power dependence of transmission in linear and circular polarizations



What limits the precision

- Even at very low density we have oscillations of the transmission intensity across the sample
- This means, that average polariton number may not be constant at given power, but for different detunings ϵ only ratio a_2/a_1 can be trusted



Measuring LPB shift

$$(T-T18 \text{ mW}) / (T+T18 \text{ mW})$$

Red → linear

Black →
circular



Blue shift, almost no broadening in
both polarizations

Negligible shift and broadening
in linear polarization



Ratio between interaction constants

Red → linear

Black →

circular . . . rson=poor
precision at zero and
strong negative detuning

$$DEL=n(a1+a2)/2$$

$$DEC=na1$$

❖ a_2 and a_1 have
different sign

❖ $|a_2|$ increases when



Tentative explanation

Different contribution to the interaction constants a_1 ($\uparrow \uparrow$) and a_2 ($\uparrow \downarrow$)

Spin independent contributions:

1) Mean field electrostatic energy (Repulsion)

1) Van-der-Waals (dipole-dipole) interaction
 $a_1 n = U_C \text{Coulomb} + U_V \text{dW}$
 (Attraction)
 $a_2 n = U_C \text{Coulomb} + U_V \text{dW}$
 $+ U_{\text{ex}} \uparrow \downarrow + U_{\text{bi}}$

$$DEC = a_1 n \quad DEL = (a_1 + a_2) n / 2$$

Spin dependent contributions:

1) Exchange interaction
 (Repulsion for $\uparrow \uparrow$ and
 Attraction for $\uparrow \downarrow$)

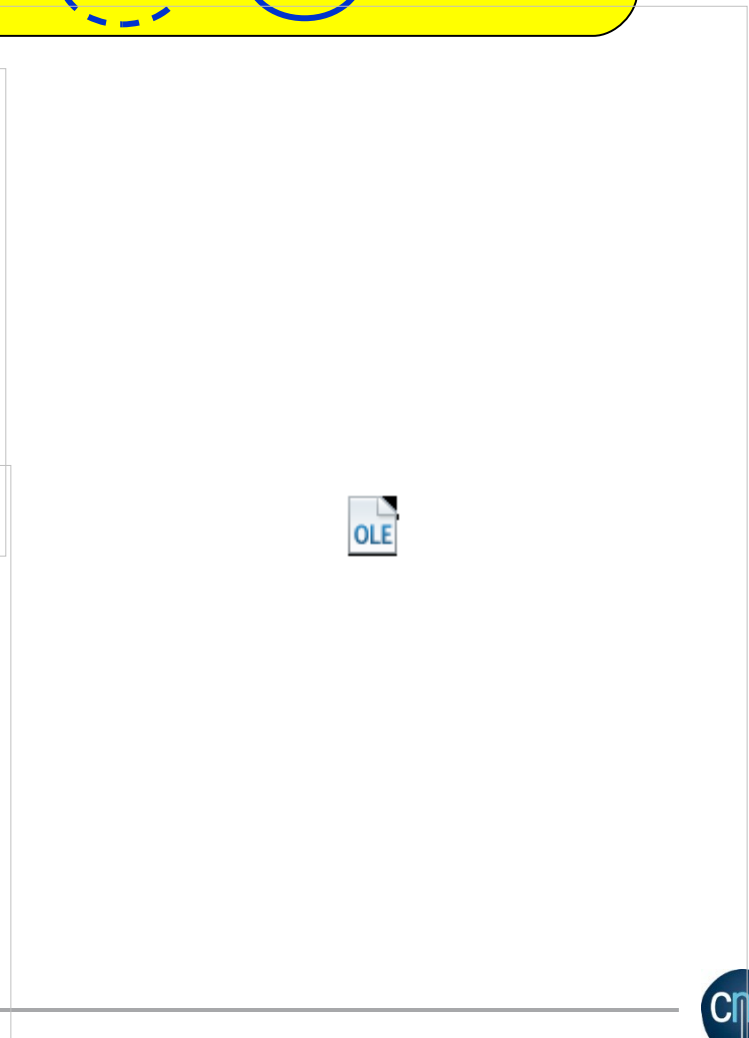
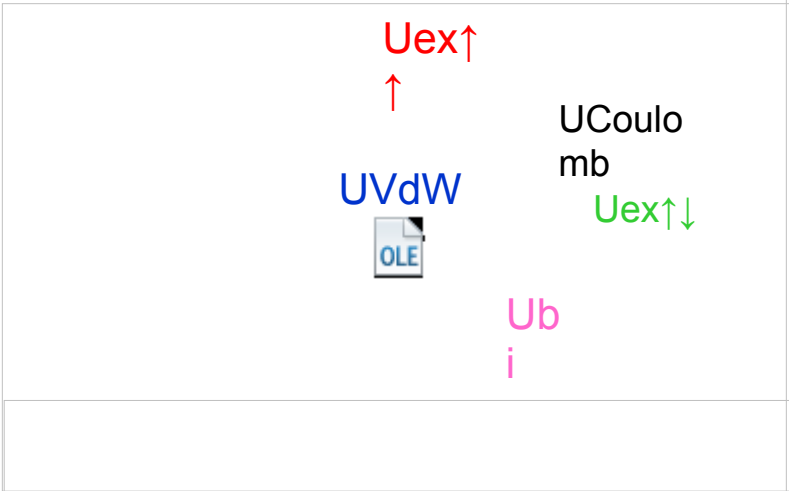
1) Bi-exciton state
 (Attraction) $\uparrow \downarrow$



measured \rightarrow $a_{1n} = \text{UCoulomb} + \text{UVdW}$ Fit of DEC

calculated \rightarrow $a_{2n} = \text{UCoulomb} + \text{UVdW} + \text{Uex}\uparrow\downarrow + \text{Ubi}$ Fit of a2/a1

$+$ $\text{Uex}\uparrow\uparrow$



Comparison with other experiments



T. Lecomte :

polarization dependent parametric
scattering in 3-coupled microcavities

P. Renucci and D. Solnyshkov:

polariton spin dynamics observed in time
and polarization resolved PL (2 μ cavities
quite similar to ours)

strong disagreement at small negative detunings!



Conclusions

I. A. Shelykh et al, SST (2010)

